Black Hole Perturbation Theory: An Introduction

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- [Regge-Wheeler](#page-24-0)
- [Zerilli](#page-42-0)

<span id="page-2-0"></span>

#### Recall

- $\rightarrow$  The spherical tensors and its orthogonality relations [\[2\]](#page-52-0):
	- Scalar  $\rightarrow$   $Y^{lm}$ :

$$
\langle Y^{lm}, Y^{l'm'} \rangle = \delta^{ll'} \delta^{mm'} \tag{1}
$$

• Vector  $\rightarrow Y_a^{lm}$  (even),  $S_a^{lm}$  (polar):

$$
\gamma^{ab}\langle S_a^{lm}, S_b^{l'm'} \rangle = l(l+1)\delta^{ll'}\delta^{mm'}
$$

$$
\gamma^{ab}\langle Y_a^{lm}, Y_b^{l'm'} \rangle = l(l+1)\delta^{ll'}\delta^{mm'}
$$

$$
\gamma^{ab}\langle Y_a^{lm}, S_b^{l'm'} \rangle = 0
$$
 (2)

• Tensor  $\rightarrow Z_{ab}^{lm}$  (even),  $S_{ab}^{lm}$  (polar):

$$
\gamma^{ac}\gamma^{bd}\langle S_{ac}^{lm}, S_{bd}^{l'm'}\rangle = l(l+1)(l+2)\delta^{ll'}\delta^{mm'}
$$

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	- Vector  $\rightarrow e_a^{e,lm}, e_a^{o,lm}$ :

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• Tensor 
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:

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\n
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 $\rightarrow$  where:

$$
h_{AB} = h_{AB}^{lm} Y^{lm} \tag{7}
$$

$$
h_{Aa} = h_A^{e,lm} e_a^{e,lm} + h_A^{o,lm} e_a^{o,lm}
$$
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$$
h_{ab} = r^2 \left( K^{lm} \gamma_{ab} Y^{lm} + G^{lm} e^{e, lm}_{ab} + 2h^{lm} e^{o, lm}_{ab} \right)
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 $\rightarrow$  Parity decomposition:

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h_{\mu\nu}(t,r,\theta,\phi) = h^e_{\mu\nu}(t,r,\theta,\phi) + h^o_{\mu\nu}(t,r,\theta,\phi),\tag{10}
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$$
h_{\mu\nu}^{e} = \left( r^{2} K^{lm} \gamma_{ab} Y^{lm} + h_{AB}^{lm} Y^{lm} + h_{A}^{e,lm} e_{a}^{e,lm} + G^{lm} e_{ab}^{e,lm} \right)_{\mu\nu}
$$

$$
h_{\mu\nu}^{o} = \left( h_{A}^{o,lm} e_{a}^{o,lm} + 2 h^{lm} e_{ab}^{o,lm} \right)_{\mu\nu}
$$
(11)

### The Regge Wheeler Gauge

 $\rightarrow$  In the Regge-Wheeler gauge:

$$
h^{odd}_{\mu\nu} = \begin{pmatrix} h^{o,lm}_{0} \\ h^{o,lm}_{1} \\ 0 \\ 0 \end{pmatrix} \begin{pmatrix} 0 & 0 & e^{o,lm}_{\theta} & e^{o,lm}_{\phi} \\ 0 & 0 & e^{o,lm}_{\theta} & e^{o,lm}_{\phi} \end{pmatrix} = \begin{pmatrix} 0 & 0 & h^{o,lm}_{0}e^{o,lm}_{\theta} & h^{o,lm}_{0}e^{o,lm}_{\phi} \\ 0 & 0 & h^{o,lm}_{1}e^{o,lm}_{\theta} & h^{o,lm}_{1}e^{o,lm}_{\phi} \\ 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \end{pmatrix}
$$

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$$

 $\rightarrow$  Symmetrizing:

$$
h^{odd}_{\mu\nu} = \begin{pmatrix} 0 & 0 & h^{o,lm}_0 \, e^{o,lm}_\theta & h^{o,lm}_0 \, e^{o,lm}_\phi \\ 0 & 0 & h^{o,lm}_1 \, e^{o,lm}_\theta & h^{o,lm}_1 \, e^{o,lm}_\phi \\ * & * & 0 & 0 \\ * & * & 0 & 0 \end{pmatrix}
$$

### The Regge-Wheeler Gauge

 $\rightarrow$  Also:

$$
h^{e,lm}_{\mu\nu}=r^2K^{lm}Y^{lm}\begin{pmatrix} 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & sin^2\theta \end{pmatrix}+Y^{lm}\begin{pmatrix} fH^{lm}_0 & H^{lm}_1 & 0 & 0 \\ H^{lm}_1 & f^{-1}H^{lm}_2 & 0 & 0 \\ 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \end{pmatrix}
$$

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$$

 $\rightarrow$  s.t.:

$$
h_{\mu\nu}^{e,lm} = \begin{pmatrix} fH_0^{lm} & H_1^{lm} & 0 & 0 \\ H_1^{lm} & f^{-1}H_2^{lm} & 0 & 0 \\ 0 & 0 & r^2K^{lm} & 0 \\ 0 & 0 & 0 & r^2\sin^2\theta K^{lm} \end{pmatrix} Y^{lm}
$$
 (12)

$$
2\delta G_{AB} = \begin{pmatrix} A_{lm}^{(0)} & A_{lm}^{(1)} \\ A_{lm}^{(1)} & A_{lm}^{(2)} \end{pmatrix} Y^{lm} \tag{13}
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$$
2\delta G_{Aa} = \alpha_A^{lm} e_a^{e,lm} + \beta_A^{lm} e_a^{o,lm} \tag{14}
$$

$$
2\delta G_{ab} = A_{lm}^{(3)} r^2 \gamma_{ab} Y^{lm} + s_{lm} e_{ab}^{e,lm} + t_{lm} e_{ab}^{o,lm} \tag{15}
$$

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$$
2\gamma^{ab}\left(\delta G_{Aa}e_b^{o,lm}\right) = \beta_A^{lm} \tag{16}
$$

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2\gamma^{ac}\gamma^{bd}\left(\delta G_{ab}e_{cd}^{o,lm}\right) = t_{lm} \tag{17}
$$

 $\rightarrow$  where:

$$
\beta_0^{lm} = f(h_{0,lm}'' + i\omega h_{1,lm}') - 2i\omega \frac{f}{r}h_1^{lm} + \left(\frac{f''}{2} + \frac{l(l+1) + f - 1}{r^2}\right)h_0^{lm}
$$
  

$$
\beta_1^{lm} = f^{-1}(i\omega h_{0,lm}' - \omega^2 h_1^{lm}) - 2i\omega \frac{f^{-1}}{r}h_0^{lm} + \left(\frac{f''}{2} + \frac{(l(l+1) - f - 1)}{r^2}\right)h_1^{lm}
$$
  

$$
t_{lm} = i\omega f^{-1}h_0^{lm} + fh_{1,lm}' + f'h_1^{lm}
$$

# <span id="page-23-0"></span>[The Master Equations](#page-23-0)

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 $\rightarrow$  Hence, there are only two independent equations:

$$
f^{-1}(i\omega h'_{0,lm} - \omega^2 h_1^{lm}) - 2i\omega \frac{f^{-1}}{r} h_0^{lm} + \left(\frac{f''}{2} + \frac{(l(l+1) - f - 1)}{r^2}\right) h_1^{lm} = 16\pi \delta T_1^{odd}
$$

$$
i\omega f^{-1} h_0^{lm} + f h'_{1,lm} + f' h_1^{lm} = 16\pi \delta T_s^{odd}
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$$

$$
i\omega f^{-1} h_0^{lm} + f h'_{1,lm} + f' h_1^{lm} = 16\pi \delta T_s^{odd}
$$

 $\rightarrow$  Making [\[1\]](#page-52-1):

$$
16\pi \delta T_1^{odd} = A
$$

$$
16\pi \delta T_s^{odd} = B
$$

$$
f^{-1}(i\omega h'_{0,lm} - \omega^2 h_1^{lm}) - 2i\omega \frac{f^{-1}}{r} h_0^{lm} + \left(\frac{f''}{2} + \frac{(l(l+1) - f - 1)}{r^2}\right) h_1^{lm} = A
$$
  

$$
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$$
  

$$
i\omega f^{-1} h_0^{lm} + f h'_{1,lm} + f' h_1^{lm} = B
$$

 $\rightarrow$  Isolating  $h_0^{lm}$  in the second equation:

$$
h_0^{lm} = \frac{if}{\omega} \left( f h'_{1,lm} + f' h_1^{lm} - B \right)
$$
 (18)

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$$
h_0^{lm} = \frac{if}{\omega} \left( f h'_{1,lm} + f' h_1^{lm} - B \right)
$$
 (18)

 $\rightarrow$  Taking the derivative on both sides:

$$
h'_{0,lm} = \left(\frac{if}{\omega} \left( f h'_{1,lm} + f' h_1^{lm} - B \right) \right)'
$$
 (19)

 $\rightarrow$  Subs eq.:

$$
h'_{0,lm} = \left(\frac{if}{\omega} \left( f h'_{1,lm} + f' h_1^{lm} - B \right) \right)'
$$
 (20)

 $\rightarrow$  into:

$$
f^{-1}(i\omega h'_{0,lm} - \omega^2 h_1^{lm}) - 2i\omega \frac{f^{-1}}{r} h_0^{lm} + \frac{(l-1)(l+2)}{r^2} h_1^{lm} = A \qquad (21)
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$$

 $\rightarrow$  We get:

$$
\left(\omega^2 + (f^2)' + \frac{1}{2}ff'' - \frac{2ff'}{r} - f\frac{(l(l+1) - f - 1)}{r^2}\right)h_1^{lm} +
$$

$$
+ h_{1,lm}''f^2 + \left(3f' - \frac{2f}{r}\right)fh_{1,lm}' = -A + (fB)' - \frac{2fB}{r}
$$
(22)

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$$
rfQ''_{lm} + rf'Q'_{lm} + \left(\frac{r\omega^2}{f} - \frac{r}{2}f'' + f' - \frac{(l(l+1) + f - 1)}{r}\right)Q_{lm} = F_{lm}(r)
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$$

 $\rightarrow$  Introducing the tortoise coordinate  $r_*(r)$  where:

$$
dr_* = \frac{dr}{f} \tag{24}
$$

 $\rightarrow$  In these coordinates, and for  $f = 1 - \frac{\sigma}{r}$  $\frac{\sigma}{r}$ :

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$$
\frac{\partial^2 Q_{lm}}{\partial r_*^2} + f \frac{\left(\frac{2\omega^2 r^2}{f} - r^2 f'' + 2r f' - 2l(l+1) - 2f + 2\right)}{2r^2} Q_{lm} = F_{lm}(r)
$$
\n(25)

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$$
\frac{\partial^2 Q_{lm}}{\partial r_*^2} + \left(\omega^2 - V_{odd}\right) Q_{lm} = F_{lm}(r) \tag{26}
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 $\rightarrow$  where:

$$
V_{odd} = \left(\frac{f''}{2} - \frac{f'}{r} - \frac{(l(l+1) + f - 1)}{r^2}\right)
$$
  

$$
F_{lm} = -16\pi \left[\delta T_1^{odd} - \delta T_s^{odd} \left(f' - \frac{2f}{r}\right) - f\left(\delta T_s^{odd}\right)'\right]
$$
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$$
(27)

 $\rightarrow$  Eq. [26](#page-38-0) is known as the **Regge-Wheeler equation**.

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 $\rightarrow$  As seen in the SageMath notebook:

$$
\frac{\partial^2 Z_{lm}}{\partial r_*^2} + \left(\omega^2 - V_{even}\right) Z_{lm} = S_{lm}(r) \tag{28}
$$

 $\rightarrow$  which is the **Zerilli equation** 

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k^{\mu}k_{\mu} = 0 \qquad k^{\mu}P_{\mu\nu} = 0 \tag{30}
$$

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 $\rightarrow$  Then we took a turn and ended up on curved spacetimes. In this context we discovered that perturbations of spherically symmetric curved spacetimes are governed by Schrodinger-like equations:

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\frac{\partial^2 Q_{lm}}{\partial r_*^2} + (\omega^2 - V_{odd}) Q_{lm} = F_{lm}(r)
$$

$$
\frac{\partial^2 Z_{lm}}{\partial r_*^2} + (\omega^2 - V_{even}) Z_{lm} = S_{lm}(r)
$$

#### That's All Folks!

Thank you!

Namarië



- <span id="page-52-1"></span>[1] Emanuele Berti. "Black Hole Perturbation Theory". In: Summer School on Gravitational-Wave Astronomy, International Center for Theoretical Sciences, Bangalore (2016). URL: <https://www.icts.res.in/event/page/3071>.
- <span id="page-52-0"></span>[2] Valeria Ferrari, Leonardo Gualtieri, and Paolo Pani. General relativity and its applications: black holes, compact stars and gravitational waves. CRC press, 2020.